

Modeling the current waveform in reactive high power impulse magnetron sputtering

Jón Tómas Guðmundsson^{1,2,3}, Daniel Lundin³,
Michael A. Raadu¹, C. Huo¹, Nils Brenning¹
and Tiberu M. Minea³

¹ Department of Space and Plasma Physics, School of Electrical Engineering,
KTH – Royal Institute of Technology, Stockholm, Sweden

² Science Institute, University of Iceland, Reykjavik, Iceland

³ Laboratoire de Physique des Gaz et Plasmas - LPGP, UMR 8578 CNRS,
Université Paris-Sud, 91405 Orsay Cedex, France

tumi@hi.is

Ruhr-Universität Bochum,
Bochum, Germany
September 29., 2016



Introduction

- Reactive sputtering, where metal targets are sputtered in a reactive gas atmosphere to deposit compound materials is of utmost importance in various technologies
- In reactive sputtering processes a reactive gas O_2 , N_2 , or CH_4 etc. is mixed to the noble working gas for oxide, nitride, or carbide deposition
- HiPIMS deposition generally gives denser, smoother films og higher crystallinity than dcMS grown films



Helmersson et al. (2006) Thin Solid Films **513** 1

Magnus et al. (2012) IEEE EDL **33** 1045

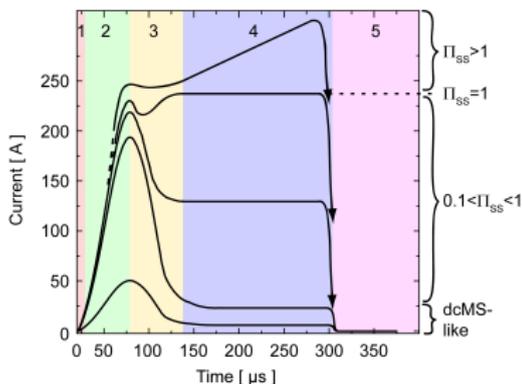


HiPIMS - Voltage - Current - Time characteristics



HiPIMS - Voltage - Current - time

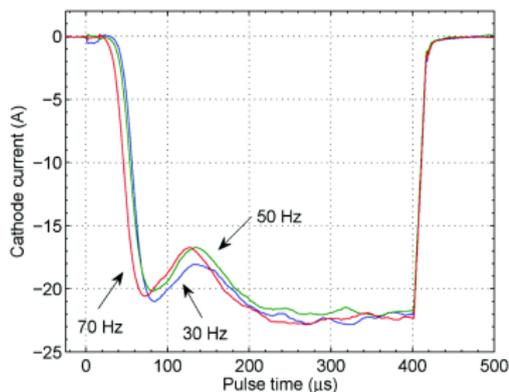
- In **non-reactive** discharge the current waveform shows an initial pressure dependent peak that is followed by a second phase that is power and material dependent
- The initial phase is dominated by gas ions, whereas the later phase has a strong contribution from self-sputtering
- The non-reactive case is well understood



From Gudmundsson et al. (2012), JVSTA **30** 030801

HiPIMS - Voltage - Current - time

- Ar discharge with Ti target
- The initial peak in current results large flux of atoms from the target
- Collisions of the sputtered atoms with the working gas result in heating and expansion of the working gas – **rarefaction**
- A significant fraction of the sputtered atoms experience electron impact ionization (the ionization mean free path ~ 1 cm) and are attracted back to the target to participate in the sputtering process – **self-sputtering**

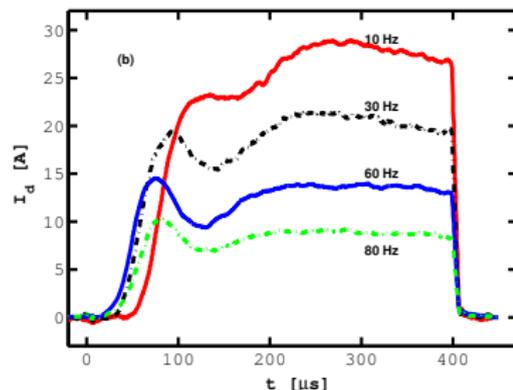


From Magnus et al. (2011) JAP **110** 083306



HiPIMS - Voltage - Current - time

- During reactive sputtering, a reactive gas is added to the inert working gas
- The current waveform in the reactive Ar/N₂ HiPIMS discharge with Ti target is highly dependent on the pulse repetition frequency
- N₂ addition changes the plasma composition and the target condition can also change due to the formation of a compound on its surface

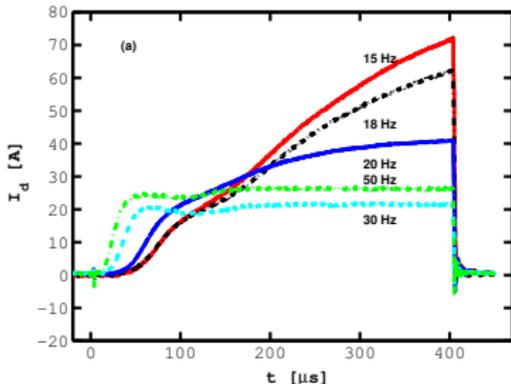


After Magnus et al. (2011) JAP **110** 083306



HiPIMS - Voltage - Current - time

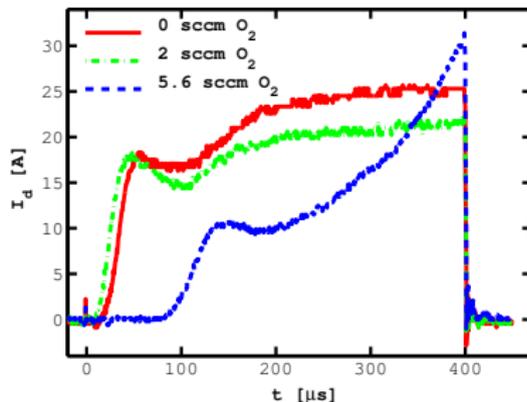
- Similarly for the Ar/O₂ discharge, the current waveform is highly dependent on the repetition frequency and applied voltage which is linked to oxide formation on the target
- The current is found to increase significantly as the frequency is lowered



After Magnus et al. (2012), JVSTA **30** 050601

HiPIMS - Voltage - Current - time

- As the oxygen flow is increased a transition to oxide mode is observed



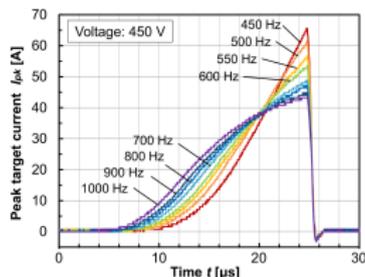
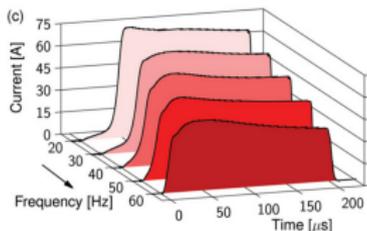
The current waveforms for an Ar/ O_2 discharge with a Ti target where the oxygen flow rate is varied – 600 V, 50 Hz and 0.6 Pa

From Gudmundsson et al. (2013), ISSP 2013, p. 192

Gudmundsson (2016) Plasma Phys. Contr. Fusion **58** 014002



HiPIMS - Voltage - Current - time



- Similar behaviour has been reported for various target and reactive gas combinations
 - The current increases with decreased repetition frequency
 - The current waveform maintains its shape for Ar/O₂ discharge with Nb target
- The current waveform becomes distinctly triangular for Ar/N₂ discharge with Hf target

From Hála et al. (2012), JPD **45** 055204

From Shimizu et al. (2016), JPD **49** 065202



HiPIMS - Voltage - Current - time

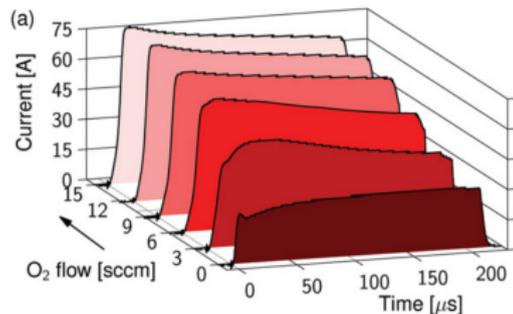
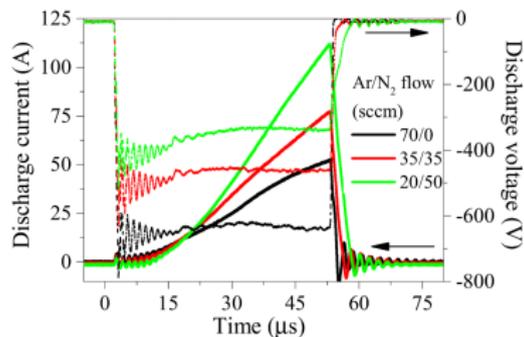
- The current increases with increased partial pressure of the reactive gas

- The current waveform becomes distinctly triangular for Ar/N₂ discharge with Al target

From Moreira et al. (2015), JVSTA **33** 021518

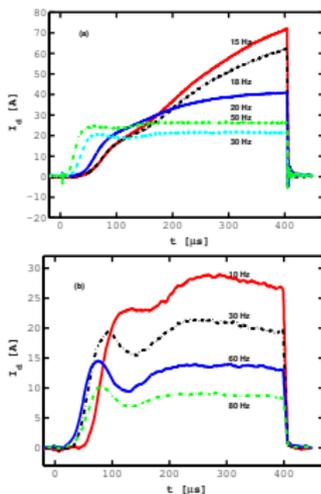
- The current waveform maintains its shape for Ar/O₂ discharge with Nb target

From Hála et al. (2012), JPD **45** 055204



HiPIMS - Voltage - Current - time

- At high frequencies, nitride or oxide is not able to form between pulses, and self-sputtering by Ti^+ -ions (singly and multiply charged) from a Ti target is the dominant process
- γ_{see} is practically zero for singly charged metal ions impacting a target of the same metal
- At low frequency, the long off-time results in a nitride or oxide layer being formed on the target surface and self-sputtering by Ti^+ - and N^+ -ions or O^+ -ions from TiN or TiO_2 takes place



From Magnus et al. (2011), JAP **110** 083306

and Magnus et al. (2012), JVSTA **30** 050601

Gudmundsson (2016) PPCF **58** 014002

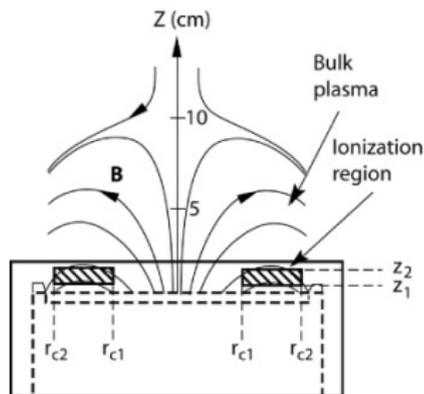


Ionization region model studies of reactive HiPIMS



Ionization region model studies of reactive HiPIMS

- The ionization region model (IRM) was developed to improve the understanding of the plasma behaviour during a HiPIMS pulse and the afterglow
- The main feature of the model is that an ionization region (IR) is defined next to the race track
- The IR is defined as an annular cylinder with outer radii r_{c2} , inner radii r_{c1} and length $L = z_2 - z_1$, extends from z_1 to z_2 axially away from the target



The definition of the volume covered by the IRM

From Raadu et al. (2011), PSST **20** 065007

Ionization region model studies of reactive HiPIMS

- The species assumed in the IRM are
 - electrons
 - argon atoms $\text{Ar}(3s^23p^6)$, warm argon atoms in the ground state Ar^W , hot argon atoms in the ground state Ar^H , Ar^m ($1s_5$ and $1s_3$) (11.6 eV), argon ions Ar^+ (15.76 eV)
 - titanium atoms $\text{Ti}(a^3F)$, titanium ions Ti^+ (6.83 eV), doubly ionized titanium ions Ti^{2+} (13.58 eV)
 - oxygen molecule in the ground state $\text{O}_2(X^3\Sigma_g^-)$, the metastable oxygen molecules $\text{O}_2(a^1\Delta_g)$ (0.98 eV) and $\text{O}_2(b^1\Sigma_g)$ (1.627 eV), the oxygen atom in the ground state $\text{O}(^3P)$, the metastable oxygen atom $\text{O}(^1D)$ (1.96 eV), the positive ions O_2^+ (12.61 eV) and O^+ (13.62 eV), and the negative ion O^-

Ionization region model studies of reactive HiPIMS

- The particle balance equation for a particular species is given by

$$\frac{dn_{\text{species}}}{dt} = \sum R_{\text{species,process}}$$

- The reaction rate $R_{j,\text{volume}}$ for a given volume reaction of a species j is

$$R_{j,\text{volume}} = k \times \prod_i n_{r,i} \quad [\text{m}^{-3}\text{s}^{-1}]$$

- The rate at which species are sputtered off the target (Ti and O) is

$$R_{n,\text{sputt}} = \frac{\sum_i \Gamma_i^{\text{RT}} S_{\text{RT}} Y_i(E)}{V_{\text{IR}}}$$

and i stands for the ion involved in the process, here $i = \text{Ti}^+, \text{Ti}^{2+}, \text{Ar}^+, \text{O}^+, \text{and } \text{O}_2^+$, Γ_i^{RT} is the flux of ion i towards the target



Ionization region model studies of reactive HiPIMS

- For each ion there is a loss rate given as

$$R_{i,\text{loss}} = - \frac{\Gamma_i^{\text{BP}} S_{\text{BP}} + \Gamma_i^{\text{RT}} S_{\text{RT}}}{V_{\text{IR}}}$$

where i stands for the particular ion and S_{BP} is the area of the annular cylinder facing the lower density plasma outside the IR (bulk plasma), and Γ_i^{BP} is the flux of ion i across the boundary towards the lower density plasma

- Thus the flux out of the IR towards the lower density plasma is reduced as required to obtain the assumed ion back-attraction probability β

$$\Gamma_i^{\text{BP}} = \left(\frac{1}{\beta} - 1 \right) \frac{S_{\text{IR}}}{S_{\text{BP}}} \Gamma_i^{\text{RT}}$$



Ionization region model studies of reactive HiPIMS

- Gas rarefaction lowers the density of the process gas inside the IR below the value in the surrounding gas reservoir, $n_{g,0}$
- This gives a back-diffusion (gain) term

$$R_{g,\text{refill}} = \frac{1}{2} v_{\text{ran}} \frac{(n_{g,0} - n_g) S_{\text{RT}}}{V_{\text{IR}}}$$

where the subscript g stands for the atoms and molecules of the working gas Ar and O₂



Ionization region model studies of reactive HiPIMS

- The return flux of recombined positive argon ions Ar^+ from the target is treated as two groups of atoms with different temperatures
 - a hot component Ar^{H} that returns from the target with a typical sputter energy of a few electron volts
 - a warm component Ar^{W} that is assumed to be embedded in the target at the ion impact, and then return to the surface and finally leave with the target temperature, at most 0.1 eV
- Thus there is a loss out of the IR

$$R_{\text{Ar}^Z, \text{loss}} = -v_{\text{ran}} n_{\text{Ar}^Z} \frac{S_{\text{RT}}}{V_{\text{IR}}}$$

where Z stands for H for hot or W for warm argon atoms



Ionization region model studies of reactive HiPIMS

- In addition there is a change in the neutral density due to kick-out by collisions with fast sputtered atoms coming from the target
- For each of the neutrals Y of the working gas, including each of the metastable states, the particle balance includes a loss term

$$R_{n,\text{kick-out}} = -\frac{1}{2} \frac{v_{\text{ran},X}}{L} F_{\text{coll}} \frac{m_X}{m_Y} \frac{\sum_i n_{X,i}}{\sum_i n_{Y,i}} n_Y$$

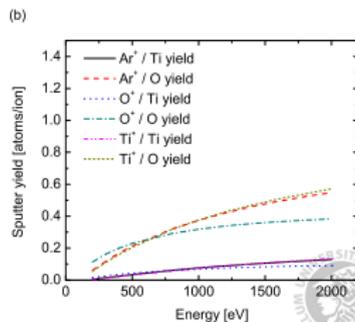
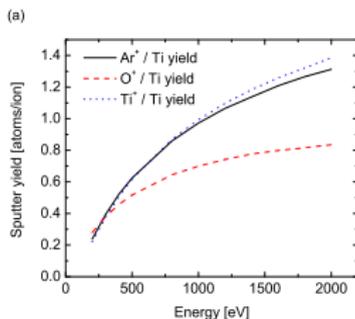
- The sum is taken over all the states of that sputtered species and

$$F_{\text{coll}} = 1 - \exp(-L/\lambda_{X,\text{gas}})$$

is the probability of a collision inside the IR, where $\lambda_{X,\text{gas}}$ is the mean free path for an atom X sputtered off the target

Ionization region model studies of reactive HiPIMS

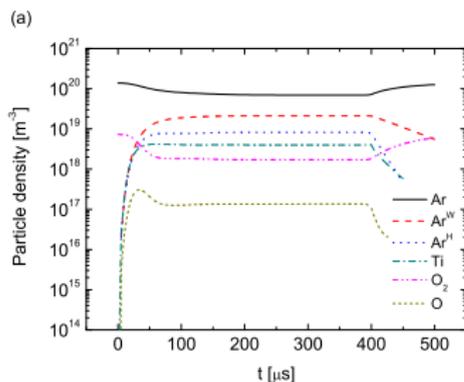
- The sputter yield for the various bombarding ions was calculated by TRIDYN for
 - **Metal mode** – Ti target
 - **Poisoned mode** – TiO₂ target
- The yields correspond to the extreme cases of either clean Ti surface and a surface completely oxidized (TiO₂ surface)
- The sputter yield is much lower for poisoned target



The sputter yield data is from Tomas Kubart, Uppsala University

Ionization region model studies of reactive HiPIMS

- The gas rarefaction is observed for the argon atoms but is more significant for the O_2 molecule
- The density of Ti atoms is higher than the O_2 density
- The atomic oxygen density of is over one order of magnitude lower than the molecular oxygen density – the dissociation fraction is low

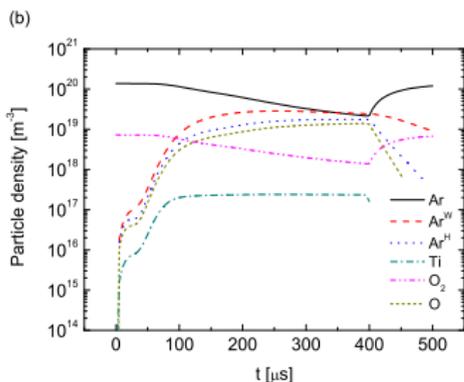


The temporal evolution of the neutral species with 5 % oxygen partial flow rate for Ar/ O_2 discharge with Ti target in **metal mode**.

Gudmundsson et al. (2016), PSST, accepted 2016

Ionization region model studies of reactive HiPIMS

- Gas rarefaction is observed for both argon atoms and O₂ molecules
- The density of Ti atoms is lower than both the O₂ density and atomic oxygen density
- The atomic oxygen density is higher than the O₂ density towards the end of the pulse



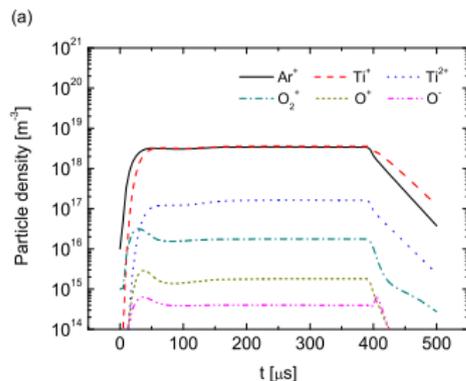
The temporal evolution of the neutral species with **5 % oxygen partial flow** rate for Ar/O₂ discharge with Ti target in **poisoned mode**.

Gudmundsson et al. (2016), PSST, accepted 2016



Ionization region model studies of reactive HiPIMS

- Ar^+ and Ti^+ -ions dominate the discharge
- Ti^{2+} -ions follow by roughly an order of magnitude lower density
- The O_2^+ and O^+ -ion density is much lower



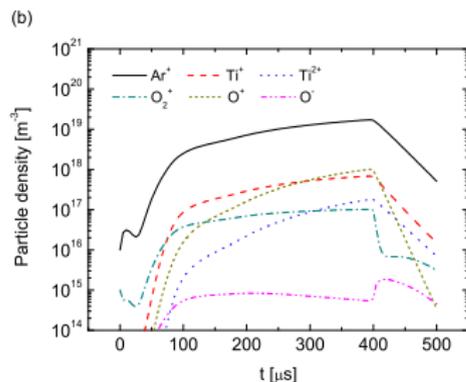
The temporal evolution of the neutral species with **5 % oxygen partial flow rate** for Ar/O_2 discharge with Ti target in **metal mode**.

Gudmundsson et al. (2016), PSST, accepted 2016



Ionization region model studies of reactive HiPIMS

- Ar^+ -ions dominate the discharge
- Ti^+ , O^+ , have very similar density, but the temporal variation is different, and the O_2^+ density is slightly lower
- The Ti^{2+} -ion density increases fast with time and overcomes the Ti^+ density towards the end of the pulse



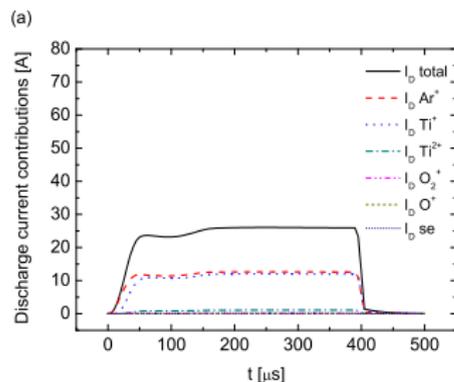
The temporal evolution of the neutral species with **5 % oxygen partial flow** rate for Ar/O_2 discharge with Ti target in **poisoned mode**.

Gudmundsson et al. (2016), PSST, accepted 2016



Ionization region model studies of reactive HiPIMS

- Ar^+ and Ti^+ -ions contribute most significantly to the discharge current



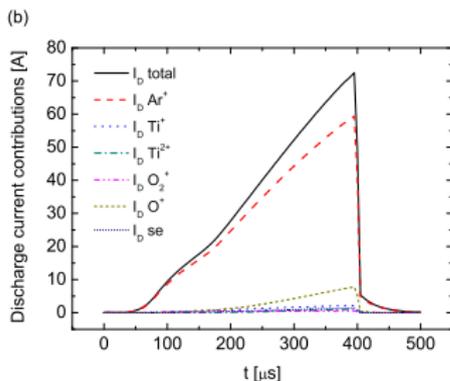
The temporal evolution of the neutral species with **5 % oxygen partial flow rate** for Ar/O_2 discharge with Ti target in **metal mode**.

Guðmundsson et al. (2016), PSST, accepted 2016



Ionization region model studies of reactive HiPIMS

- Ar^+ contribute most significantly to the discharge current – almost solely
- The contribution of secondary electron emission is very small



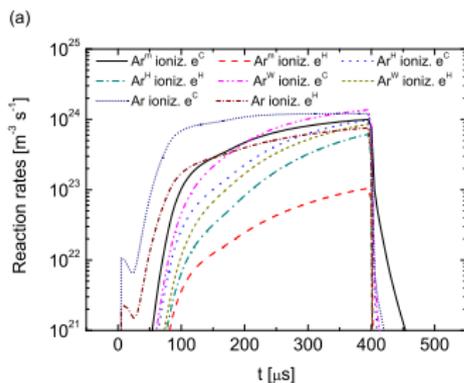
The temporal evolution of the neutral species with **5 % oxygen partial flow rate** for Ar/O_2 discharge with Ti target in **poisoned mode**.

Gudmundsson et al. (2016), PSST, accepted 2016



Ionization region model studies of reactive HiPIMS

- Recycling of ionized atoms coming from the target are required for the current generation in both modes of operation
- In the metal mode self-sputter recycling dominates and in the poisoned mode working gas recycling dominates
- The dominating type of recycling determines the discharge current waveform



The temporal variations of the reaction rates for electron impact ionization of the argon atoms (ground state plus metastable) in poisoned mode.

Gudmundsson et al. (2016), PSST, accepted 2016



Ionization region model studies of reactive HiPIMS

- In the metal mode sheath energization was found to be only 10 %
 - same range as the results reported earlier for an Al target
- For the poisoned mode the sheath energization was 30 %, with a rising trend, at the end of the pulse
- This is due to the secondary electron emission
 - In the poisoned mode essentially all the ions (mainly Ar^+ , but also O^+ and Ti^{2+} towards the end of the pulse) contribute to the secondary electron emission
 - In the metal mode only half of the ions contribute to the secondary electron emission (Ar^+) while the other half does not contribute at all ($\gamma_{\text{Ti}^+} = 0.0$)



Summary



Summary

- An ionization region model was used to explore the plasma composition during the high power pulse
- Comparison was made between the metal mode and the poisoned mode
 - In metal mode Ar^+ and Ti^+ -ions dominate the discharge and are of the same order of magnitude
 - In poisoned mode Ar^+ -ions dominate the discharge and two orders of magnitude lower, Ti^+ , O^+ , have very similar density, with the O_2^+ density slightly lower
 - In the metal mode Ar^+ and Ti^+ -ions contribute most significantly to the discharge current while in poisoned mode Ar^+ dominate
- In the metal mode self-sputter recycling dominates and in the poisoned mode working gas recycling dominates – the dominating type of recycling determines the discharge current waveform



The slides can be downloaded at

<http://langmuir.raunvis.hi.is/~tumi/ranns.html>

- The experimental work was made in collaboration with
 - Dr. Fridrik Magnus, Uppsala University, Uppsala, Sweden
 - Tryggvi K. Tryggvason, University of Iceland
- We got help with the sputtering yields from
 - Dr. Tomas Kubart, Uppsala University, Uppsala, Sweden
- and the project is funded by
 - Icelandic Research Fund Grant No. 130029-053
 - Swedish Government Agency for Innovation Systems (VINNOVA) contract no. 2014-04876,



References

- Aiempanakit, M., A. Aijaz, D. Lundin, U. Helmersson, and T. Kubart (2013). Understanding the discharge current behavior in reactive high power impulse magnetron sputtering of oxides. *Journal of Applied Physics* 113(13), 133302.
- Gudmundsson, J. T. (2016). On reactive high power impulse magnetron sputtering. *Plasma Physics and Controlled Fusion* 58(1), 014002.
- Gudmundsson, J. T., N. Brenning, D. Lundin, and U. Helmersson (2012). The high power impulse magnetron sputtering discharge. *Journal of Vacuum Science and Technology A* 30(3), 030801.
- Gudmundsson, J. T., F. Magnus, T. K. Tryggvason, S. Shayestehaminzadeh, O. B. Sveinsson, and S. Olafsson (2013). Reactive high power impulse magnetron sputtering. In *Proceedings of the XII International Symposium on Sputtering and Plasma Processes (ISSP 2013)*, pp. 192–194.
- Gudmundsson, J. T., D. Lundin, N. Brenning, M. A. Raadu, C. Huo, and T. M. Minea (accepted 2016). An ionization region model of the reactive Ar/O₂ high power impulse magnetron sputtering discharge. *Plasma Sources Science and Technology*.
- Hála, M., J. Čapek, O. Zabeida, J. E. Klemberg-Sapieha, and L. Martinu (2012). Hysteresis - free deposition of niobium oxide films by HiPIMS using different pulse management strategies. *Journal of Physics D: Applied Physics* 45(5), 055204.



References

- Helmersson, U., M. Lattemann, J. Bohlmark, A. P. Ehasarian, and J. T. Gudmundsson (2006). Ionized physical vapor deposition (IPVD): A review of technology and applications. *Thin Solid Films* 513(1-2), 1–24.
- Magnus, F., A. S. Ingason, S. Olafsson, and J. T. Gudmundsson (2012). Nucleation and resistivity of ultrathin TiN films grown by high power impulse magnetron sputtering. *IEEE Electron Device Letters* 33(7), 1045 – 1047.
- Magnus, F., O. B. Sveinsson, S. Olafsson, and J. T. Gudmundsson (2011). Current-voltage-time characteristics of the reactive Ar/N₂ high power impulse magnetron sputtering discharge. *Journal of Applied Physics* 110(8), 083306.
- Magnus, F., T. K. Tryggvason, S. Olafsson, and J. T. Gudmundsson (2012). Current-voltage-time characteristics of the reactive Ar/O₂ high power impulse magnetron sputtering discharge. *Journal of Vacuum Science and Technology A* 30(5), 050601.
- Moreira, M. A., T. Törndahl, I. Katardjiev, and T. Kubart (2015). Deposition of highly textured AlN thin films by reactive high power impulse magnetron sputtering. *Journal of Vacuum Science and Technology A* 33(2), 021518.
- Raadu, M. A., I. Axnäs, J. T. Gudmundsson, C. Huo, and N. Brenning (2011). An ionization region model for high power impulse magnetron sputtering discharges. *Plasma Sources Science and Technology* 20(6), 065007.
- Shimizu, T., M. Villamayor, D. Lundin, and U. Helmersson (2015). Process stabilization by peak current regulation in reactive high-power impulse magnetron sputtering of hafnium nitride. *Journal of Physics D: Applied Physics* 49(6), 065202.
- Toneli, D. A., R. S. Pessoa, M. Roberto, and J. T. Gudmundsson (2015). On the formation and annihilation of the singlet molecular metastables in an oxygen discharge. *Journal of Physics D: Applied Physics* 48(32), 325202.

