On recycling in high power impulse magnetron sputtering discharges

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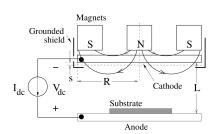






Introduction

- Magnetron sputtering has been a highly successfull technique that is essential in a number of industrial applications
- A magnet is placed at the back of the cathode target with the pole pieces at the center and perimeter
- The magnetic field confines the energetic electrons near the cathode
- The electrons undergo numerous ionizing collisions before being lost to a grounded surface





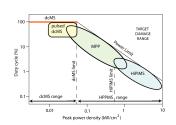






High power impulse magnetron sputtering discharge

- High ionization of sputtered material requires very high density plasma
- In a conventional dc magnetron sputtering discharge the power density (plasma density) is limited by the thermal load on the target
- High power pulsed magnetron sputtering (HPPMS)
- In a HiPIMS discharge a high power pulse is supplied for a short period
 - low frequency
 - low duty cycle
 - low average power



Gudmundsson et al. (2012) JVSTA 30 030801

Power density limits
p_t = 0.05 kW/cm² dcMS limit
p_t = 0.5 kW/cm² HiPIMS limit

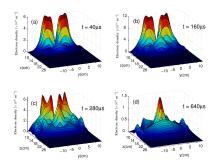








High power impulse magnetron sputtering discharge



(After Bohlmark et al. (2005), IEEE Trans. Plasma Sci. 33 346)

- Temporal and spatial variation of the electron density
- Ar discharge at 20 mTorr, Ti target, pulse length 100 μ s
- The electron density in the substrate vicinity is of the order of $10^{18} 10^{19}$ m⁻³ ionization mean free path $\lambda_{iz} \sim 10^{19}$



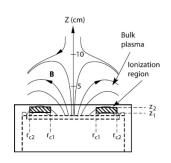


Ionization region model studies of HiPIMS discharges





- The ionization region model (IRM) was developed to improve the understanding of the plasma behaviour during a HiPIMS pulse and the afterglow
- The main feature of the model is that an ionization region (IR) is defined next to the race track
- The IR is defined as an annular cylinder with outer radii r_{c2} , inner radii r_{c1} and length $L = z_2 z_1$, extends from z_1 to z_2 axially away from the target



The definition of the volume covered by the IRM From Raadu et al. (2011) PSST **20** 065007







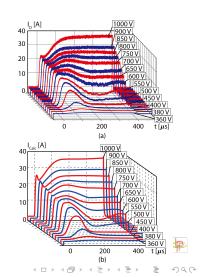


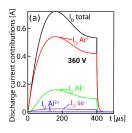
- The temporal development is defined by a set of ordinary differential equations giving the first time derivatives of
 - the electron energy
 - the particle densities for all the particles
- The species assumed in the of-IRM are
 - cold electrons e^C (Maxwellian), hot electrons e^H (sheath acceleration)
 - argon atoms Ar(3s²3p⁶), warm argon atoms in the ground state Ar^W, hot argon atoms in the ground state Ar^H, Ar^m (1s₅ and 1s₃) (11.6 eV), argon ions Ar⁺ (15.76 eV)
 - titanium atoms Ti(a³F), titanium ions Ti⁺ (6.83 eV), doubly ionized titanium ions Ti²⁺ (13.58 eV)
 - aluminium atoms Al(²P_{1/2}), aluminium ions Al⁺ (5.99 eV), doubly ionized aluminium ions Al²⁺ (18.8 eV)

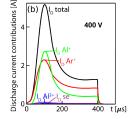
- The model is constrained by experimental data input and fitted to reproduce the measured discharge current and voltage curves, I_D(t) and V_D(t), respectively
- Two model fitting parameters were found to be sufficient for a discharge with Al target
 - V_{IR} accounts for the power transfer to the electrons
 - β is the probability of back-attraction of ions to the target

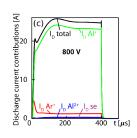
From Huo et al. (2017) JPD 50 354003

Experimental data from Anders et al. (2007) JAP 102 113303









- A non-reactive discharge with 50 mm diameter Al target
- Current composition at the target surface

From Huo et al. (2017) JPD 50 354003

Experimental data from Anders et al. (2007) JAP 10



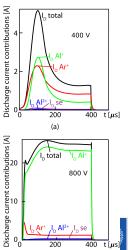




- When the discharge is operated at 400 V the contributions of Al⁺ and Ar⁺-ions to the discharge current are very similar
- At 800 V Al⁺-ions dominate the discharge current (self-sputtering) while the contribution of Ar⁺ is below 10 % except at the initiation of the pulse

From Huo et al. (2017) JPD 50 354003

Experimental data from Anders et al. (2007) JAP 102 113303





- A primary current I_{prim} is defined as ions of the working gas, here Ar⁺, that are ionized for the first time and then drawn to the target
- This is the dominating current in dc magnetron sputtering discharges
- This current has a critical upper limit

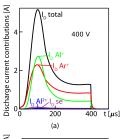
$$I_{\mathrm{crit}} = S_{\mathrm{RT}} e
ho_{\mathrm{g}} \sqrt{rac{1}{2\pi m_{\mathrm{g}} k_{\mathrm{B}} T_{\mathrm{g}}}} = S_{\mathrm{RT}} e n_{\mathrm{g}} \sqrt{rac{k_{\mathrm{B}} T_{\mathrm{g}}}{2\pi m_{\mathrm{g}}}}$$

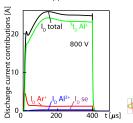
 Discharge currents I_D above I_{crit} are only possible if there is some kind of recycling of atoms that leave the target, become subsequently ionized and then are drawn back to the target

- For the 50 mm diameter Al target the critical current is $I_{\rm crit} \approx 7$ A
- The experiment is operated from far below I_{crit} to high above it, up to 36 A.
- With increasing current I_{prim} gradually becomes a very small fraction of the total discharge current I_D
- The current becomes mainly carried by singly charged Al⁺-ions, meaning that self-sputter recycling or the current I_{SS-recycle} dominates

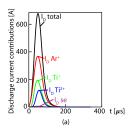
From Huo et al. (2017) JPD 50 354003

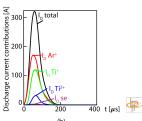
Experimental data from Anders et al. (2007) JAP 102 113303





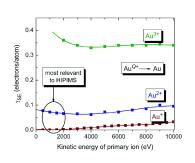
- For discharges with Ti target the peak current is far above the critical current (up to 650 A, while $I_{\rm crit} \approx$ 19 A)
- However, this discharge shows close to a 50/50 combination of self-sputter recycling I_{SS-recycle} and working gas-recycling I_{gas-recycle}
- Almost 2/3 of the current to the target is here carried by Ar⁺ and Ti²⁺-ions, which both can emit secondary electrons upon target bombardment, and this gives a significant sheath energization





From Huo et al. (2017) JPD 50 354003

- Recall that singly charged metal ions cannot create the secondary electrons for metal self-sputtering (γ_{SE} is practically zero)
- The first ionization energies of many metals are insufficient to overcome the workfunction of the target material
- For the discharge with Al target operated at high voltage, self-sputter dominated, the effective secondary electron emission is essentially zero



From Anders (2008) APL 92 201501

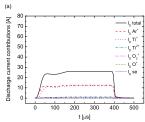


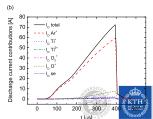






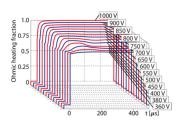
- Reactive HiPIMS
- Ar/O₂ discharge with Ti target
- For this system $I_{\rm crit} \approx 5$ A
- In the metal mode Ar⁺ and Ti⁺-ions contribute roughly equally to the current – combined self-sputter recycling and working gas recycling
- In the poisoned mode the current increaes and Ar⁺-ions dominate the current – working gas recycling







- For the Al target, Ohmic heating is in the range of 87 % (360 V) to 99 % (1000 V)
- The domination of Al⁺-ions, which have zero secondary electron emission yield, has the consequence that there is negligible sheath energization
- The ionization threshold for twice ionized Al²⁺, 18.8 eV, is so high that few such ions are produced



From Huo et al. (2017) JPD 50 354003

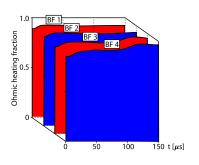








- For a Ti target Ohmic heating is about 92 %
 - Both Ar⁺ and Ti²⁺-ions contribute to creation of secondary electrons
- For Ti target in Ar/O₂ mixture
 - In the metal mode Ohmic heating is found to be 90 % during the plateau phase of the discharge pulse
 - For the poisoned mode Ohmic heating is 70 % with a decreasing trend, at the end of the pulse



From Huo et al. (2017) JPD 50 354003

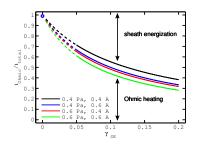








- Ohmic heating is also very significant in dc magnetron sputtering discharges
- The relative contributions to the total ionization ι_{total} due to Ohmic heating, ι_{Ohmic} , and sheath energization, ι_{sheath}
- A blue circle marks the HiPIMS study modelled by Huo et al. (2013)
- Note that this HiPIMS case $\gamma_{\rm SE,eff}$ is consistent with the dcMS cases



The generalized recycling model





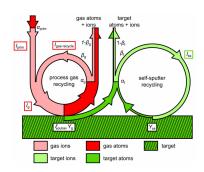
A working gas-sputtering parameter

$$\pi_{\rm g} = \alpha_{\rm g} \beta_{\rm g} \xi_{\rm pulse}$$

where

- $\alpha_{\rm g}$ is ionization probability
- β_g is back attraction probability
- $\xi_{\text{pulse}} = 1$ is return fraction in a pulse
- The total current carried by working gas ions

$$I_{g} = I_{prim} + I_{gas-recycle} = I_{prim} \left(1 + \frac{\pi_{g}}{1 - \pi_{g}} \right)$$





The total self-sputter current is

$$\textit{I}_{SS} = \textit{I}_{g} \left(\frac{\textit{Y}_{g}}{\textit{Y}_{SS}} \frac{\pi_{SS}}{1 - \pi_{SS}} \right)$$

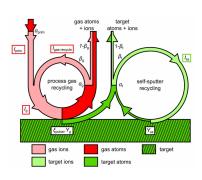
where the self-sputter parameter is

$$\pi_{\rm SS} = \alpha_{\rm t} \beta_{\rm t} \, Y_{\rm SS}$$

The total discharge current is

$$I_{D} = I_{\text{prim}} + I_{\text{gas-recycle}} + I_{\text{SS}}$$

$$= I_{\text{prim}} \left(1 + \frac{\pi_{\text{g}}}{1 - \pi_{\text{g}}} \right) \left(1 + \frac{Y_{\text{g}}}{Y_{\text{SS}}} \frac{\pi_{\text{SS}}}{1 - \pi_{\text{SS}}} \right)$$



From Brenning et al. (2017) PSST 26 125003.





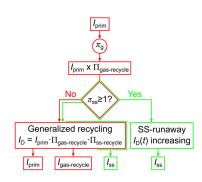




The discharge current

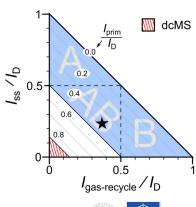
$$I_{\rm D} = I_{\rm prim} \Pi_{\rm gas-recycle} \Pi_{\rm SS-recycl}$$

- I_{prim} is the seed current acts as a seed to the whole discharge current and has an upper limit I_{crit}
- I_{prim}Π_{gas-recycle} is the seed current for the self-sputter process
- If $\pi_{\rm SS} >$ 1 the discharge goes into SS-runaway





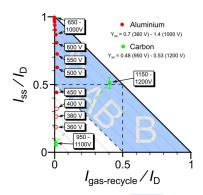
- Recycling map
- A graph in which the ion current mix of I_{prim}, I_{gas-recycle}, and I_{SS} to the target in a magnetron discharge is defined by a point
- The value of $I_{\text{prim}}/I_{\text{D}} = 39$ %, can be read on the diagonal lines ($Y_{\text{SS}} = 0.5$)
- $I_{\rm prim}/I_{\rm D} \geq$ 0.85 defines the dcMS regime
- For $I_{\rm SS}/I_{\rm D}>0.5$ we have the SS-recycle dominated range A
- For I_{gas-recycle}/I_D > 0.5 we have the gas-recycle dominated range B







- The discharge with AI target moves from the dcMS regime to the HiPIMS discharge regime with increased discharge voltage – type A
- A discharge with carbon target jumps from the dcMS regime to the HiPIMS regime – both SS recycling and working gas recycling play a role – intermediate type AB



 Recycling map for five different targets with varying self-sputter yield

•
$$Cu - Y_{SS} = 2.6$$

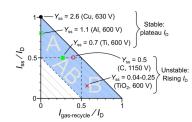
• Al –
$$Y_{SS} = 1.1$$

• Ti –
$$Y_{SS} = 0.7$$

$$\circ$$
 C – $Y_{SS} = 0.5$

$$\bullet$$
 TiO₂ - $Y_{SS} = 0.04 - 0.25$

- For very high self-sputter yields
 Y_{SS} > 1, the discharges above I_{crit} are of type A with dominating SS-recycling
- For very low self-sputter yields $Y_{\rm SS} < 0.2$, the discharges above $I_{\rm crit}$ are of **type B** with dominating **working gas recycling**



From Brenning et al. (2017),
PSST **26** 125003

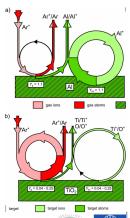






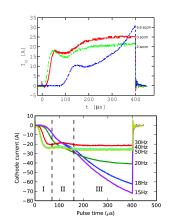


- Recycling loops
- Discharge with Al target SS recycling dominates
 - high self sputter yield
- Reactive discharge with TiO₂ target working gas recycling dominates
 - low self sputter yield



HiPIMS - Voltage - Current - time

- For Ar/O₂ discharge with Ti target
- At high frequencies, oxide is not able to form between pulses, and self-sputtering recycling by Ti⁺-ions is the dominant process
- At low frequency, the long off-time results in an oxide layer being formed (TiO₂) on the target surface and working gas recycling dominates – triangular current waveform



Summary

- For high currents the discharge with Al target develops almost pure self-sputter recycling, while the discharge with Ti target exhibits close to a 50/50 combination of self-sputter recycling and working gas-recycling
- For very high self-sputter yields, above approximately $Y_{\rm SS} \approx$ 1, the discharges above $I_{\rm crit}$ are of type A with
 - dominating SS-recycling
 - very little secondary electron emission
 - little sheath energization of electrons
- For very low self-sputter yields, below approximately $Y_{\rm SS} \approx 0.2$, the discharges above $I_{\rm crit}$ are of type B with
 - dominating working gas recycling
 - significant secondary electron emission
 - significant sheath energization of electrons.
- The fraction of the total electron heating that is attributable to Ohmic heating is over 90 % in the HiPIMS discharge





Thank you for your attention

The slides can be downloaded at http://langmuir.raunvis.hi.is/~tumi/ranns.html and the project is funded by

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