The role of recycling in pulsed sputtering magnetrons

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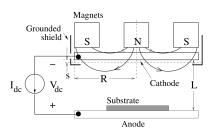
> 70th Annual Gaseous Electronics Conference Pittsburgh, Pennsylvania November 9., 2017





Introduction

- Magnetron sputtering has been a highly successfull technique that has a number of industrial applications
- A magnet is placed at the back of the cathode target with the pole pieces at the center and perimeter
- The magnetic field confines the energetic electrons near the cathode
- The electrons undergo numerous ionizing collisions before being lost to a grounded surface



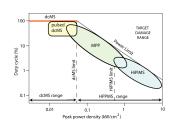






High power impulse magnetron sputtering discharge

- High ionization of sputtered material requires very high density plasma
- In a conventional dc magnetron sputtering discharge the power density (plasma density) is limited by the thermal load on the target
- High power pulsed magnetron sputtering (HPPMS)
- In a HiPIMS discharge a high power pulse is supplied for a short period
 - low frequency
 - low duty cycle
 - low average power



Gudmundsson et al. (2012), JVSTA 30 030801

Power density limits $p_t = 0.05 \text{ kW/cm}^2 \text{ dcMS limit}$ $p_t = 0.5 \text{ kW/cm}^2 \text{ HiPIMS limit}$





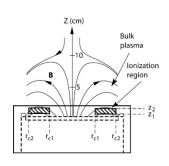


Ionization region model studies of HiPIMS discharges





- The ionization region model (IRM) was developed to improve the understanding of the plasma behaviour during a HiPIMS pulse and the afterglow
- The main feature of the model is that an ionization region (IR) is defined next to the race track
- The IR is defined as an annular cylinder with outer radii r_{c2}, inner radii r_{c1} and length
 L = z₂ z₁, extends from z₁ to z₂ axially away from the target



The definition of the volume covered by the IRM From Raadu et al. (2011), PSST **20** 065007



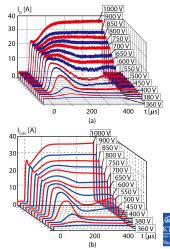


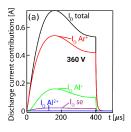


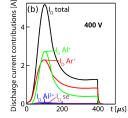
- The temporal development is defined by a set of ordinary differential equations giving the first time derivatives of
 - the electron energy
 - the particle densities for all the particles
- The species assumed in the of-IRM are
 - cold electrons e^C (Maxwellian), hot electrons e^H (sheath acceleration)
 - argon atoms Ar(3s²3p⁶), warm argon atoms in the ground state Ar^W, hot argon atoms in the ground state Ar^H, Ar^m (1s₅ and 1s₃) (11.6 eV), argon ions Ar⁺ (15.76 eV)
 - titanium atoms Ti(a³F), titanium ions Ti⁺ (6.83 eV), doubly ionized titanium ions Ti²⁺ (13.58 eV)
 - aluminium atoms Al(²P_{1/2}), aluminium ions Al⁺ (5.99 eV), doubly ionized aluminium ions Al²⁺ (18.8 eV)

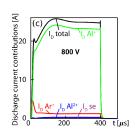
- The model is constrained by experimental data input and fitted to reproduce the measured discharge current and voltage curves, I_D(t) and V_D(t), respectively
- Two model fitting parameters were found to be sufficient for a discharge with Al target
 - V_{IR} accounts for the power transfer to the electrons
 - β is the probability of back-attraction of ions to the target

From Huo et al. (2017), JPD 50 354003









- A non-reactive discharge with Al target
- Current composition at the target surface

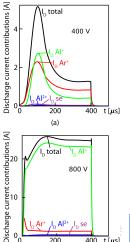
From Huo et al. (2017), JPD 50 354003





- When the discharge is operated at 400 V the contributions of Al⁺ and Ar⁺-ions to the discharge current are very similar
- At 800 V Al⁺-ions dominate the discharge current (self-sputtering) while the contribution of Ar⁺ is below 10 % except at the initiation of the pulse

From Huo et al. (2017), JPD 50 354003







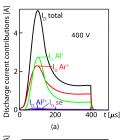
- A primary current I_{prim} is defined as ions of the working gas, here Ar⁺, that are ionized for the first time and then drawn to the target
- This is the dominating current in dc magnetron sputtering discharges
- This current has a critical upper limit

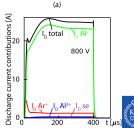
$$I_{\mathrm{crit}} = S_{\mathrm{RT}} e
ho_{\mathrm{g}} \sqrt{rac{1}{2\pi m_{\mathrm{g}} k_{\mathrm{B}} T_{\mathrm{g}}}} = S_{\mathrm{RT}} e n_{\mathrm{g}} \sqrt{rac{k_{\mathrm{B}} T_{\mathrm{g}}}{2\pi m_{\mathrm{g}}}}$$

 Discharge currents I_D above I_{crit} are only possible if there is some kind of recycling of atoms that leave the target, become subsequently ionized and then are drawn back to the target

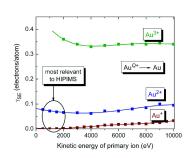
- For the AI target the critical current is $I_{\rm crit} \approx 7~{\rm A}$
- The experiment is operated from far below I_{crit} to high above it, up to 36 A.
- With increasing current I_{prim} gradually becomes a very small fraction of the total discharge current I_D
- The current becomes mainly carried by singly charged Al⁺-ions, meaning that self-sputter recycling or the current I_{SS-recycle} dominates

From Huo et al. (2017), JPD 50 354003





- Recall that singly charged metal ions cannot create the secondary electrons – for metal self-sputtering (γ_{SE} is practically zero)
- The first ionization energies of many metals are insufficient to overcome the workfunction of the target material
- For the discharge with Al target operated at high voltage, self-sputter dominated, the effective secondary electron emission is essentially zero



From Anders (2008), APL 92 201501

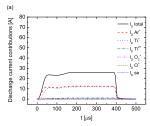


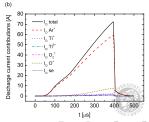




- Reactive HiPIMS
- Ar/O₂ discharge with Ti target
- In the metal mode Ar⁺ and Ti⁺-ions contribute roughly equally to the current – combined self-sputter recycling and working gas recycling
- In the poisoned mode the current increaes and Ar⁺-ions dominate the current – working gas recycling

From Gudmundsson et al. (2016), PSST 25(6) 065004







The generalized recycling model





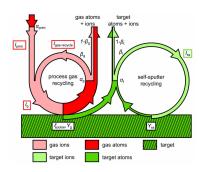
A working gas-sputtering parameter

$$\pi_{\rm g} = \alpha_{\rm g} \beta_{\rm g} \xi_{\rm pulse}$$

where

- $\alpha_{\rm g}$ is ionization probability
- β_g is back attraction probability
- $\xi_{\text{pulse}} = 1$ is return fraction in a pulse
- The total current carried by working gas ions

$$I_{\rm g} = I_{\rm prim} + I_{\rm gas-recycle} = I_{\rm prim} \left(1 + \frac{\pi_{\rm g}}{1 - \pi_{\rm g}} \right)$$



From Brenning et al. (2017), PSST 26 1250

The total self-sputter current is

$$\textit{I}_{SS} = \textit{I}_{g} \left(\frac{\textit{Y}_{g}}{\textit{Y}_{SS}} \frac{\pi_{SS}}{1 - \pi_{SS}} \right)$$

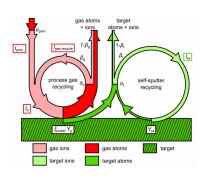
where the self-sputter parameter is

$$\pi_{\rm SS} = \alpha_{\rm t} \beta_{\rm t} \, Y_{\rm SS}$$

The total discharge current is

$$I_{D} = I_{\text{prim}} + I_{\text{gas-recycle}} + I_{\text{SS}}$$

$$= I_{\text{prim}} \left(1 + \frac{\pi_{\text{g}}}{1 - \pi_{\text{g}}} \right) \left(1 + \frac{Y_{\text{g}}}{Y_{\text{SS}}} \frac{\pi_{\text{SS}}}{1 - \pi_{\text{SS}}} \right)$$



From Brenning et al. (2017), PSST 26 125003.



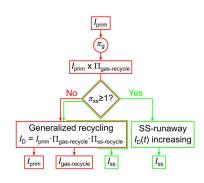




The discharge current

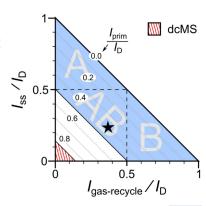
$$I_{\rm D} = I_{\rm prim} \Pi_{\rm gas-recycle} \Pi_{\rm SS-recycl}$$

- I_{prim} is the seed current acts as a seed to the whole discharge current and has an upper limit I_{crit}
- I_{prim}Π_{gas-recycle} is the seed current for the self-sputter process
- If $\pi_{\rm SS} >$ 1 the discharge goes into SS-runaway



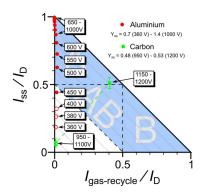
From Brenning et al. (2017), PSST 26

- Recycling map
- A graph in which the ion current mix of I_{prim}, I_{gas-recycle}, and I_{SS} to the target in a magnetron discharge is defined by a point
- The value of $I_{\text{prim}}/I_{\text{D}}=39$ %, can be read on the diagonal lines ($Y_{\text{SS}}=0.5$)
- $I_{\rm prim}/I_{\rm D}=0.85$ defines the dcMS regime
- For $I_{\rm SS}/I_{\rm D}>0.5$ we have the SS-recycle dominated range A
- For $I_{\rm gas-recycle}/I_{\rm D}>0.5$ we have the gas-recycle dominated range B





- The discharge with Al target moves from the dcMS regime to the HiPIMS discharge regime with increased discharge voltage – type A
- A discharge with carbon target jumps from the dcMS regime to the HiPIMS regime – both SS recycling and working gas recycling play a role – intermediate type AB



 Recycling map for five different targets with varying self-sputter yield

•
$$Cu - Y_{SS} = 2.6$$

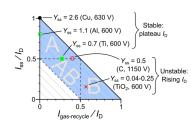
• Al –
$$Y_{SS} = 1.1$$

• Ti –
$$Y_{SS} = 0.7$$

•
$$C - Y_{SS} = 0.5$$

•
$$TiO_2 - Y_{SS} = 0.04 - 0.25$$

- For very high self-sputter yields
 Y_{SS} > 1, the discharges above I_{crit} are of type A with dominating SS-recycling
- For very low self-sputter yields
 Y_{SS} < 0.2, the discharges above I_{crit} are
 of type B with dominating working gas
 recycling



From Brenning et al. (2017),
PSST **26** 125003







- For high currents the discharge with Al target develops almost pure self-sputter recycling, while the discharge with Ti target exhibits close to a 50/50 combination of self-sputter recycling and working gas-recycling
- For very high self-sputter yields, above approximately $Y_{\rm SS} \approx$ 1, the discharges above $I_{\rm crit}$ are of type A with
 - dominating SS-recycling
 - very little secondary electron emission
 - little sheath energization of electrons
- For very low self-sputter yields, below approximately $Y_{\rm SS} \approx$ 0.2, the discharges above $I_{\rm crit}$ are of type B with
 - dominating process gas recycling
 - significant secondary electron emission
 - significant sheath energization of electrons.
- The fraction of the total electron heating that is attributable to Ohmic heating is over 90 % in the HiPIMS discharge



Thank you for your attention

The slides can be downloaded at http://langmuir.raunvis.hi.is/~tumi/ranns.html and the project is funded by

- Icelandic Research Fund Grant No. 130029
- Swedish Government Agency for Innovation Systems (VINNOVA) contract no. 2014-04876





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