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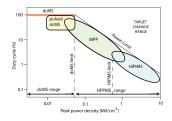
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8th International Conference on Fundamentals and Industrial Applications of HiPIMS, Braunschweig, Germany, June 14., 2017



Introduction

- High ionization of sputtered material requires very high density plasma
- In a conventional dc magnetron sputtering discharge the power density (plasma density) is limited by the thermal load on the target
- High power pulsed magnetron sputtering (HPPMS)
- In a HiPIMS discharge a high power pulse is supplied for a short period
 - Iow frequency
 - Iow duty cycle
 - Iow average power



Gudmundsson et al. (2012), JVSTA 30 030801

Power density limits

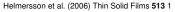
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- $p_{\rm t}=0.05~{\rm kW/cm^2}~{\rm dcMS}~{\rm limit}$
- $p_t = 0.5 \text{ kW/cm}^2 \text{ HiPIMS limit}$



Introduction

- Reactive sputtering, where metal targets are sputtered in a reactive gas atmosphere to deposit compound materials is of utmost importance in various technologies
- In reactive sputtering processes a reactive gas O₂, N₂, or CH₄ etc. is mixed to the noble working gas for oxide, nitride, or carbide deposition
- HiPIMS deposition generally gives denser, smoother films and higher crystallinity than dcMS grown films



Magnus et al. (2012) IEEE EDL 33 1045





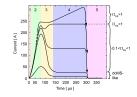
Voltage - Current - Time characteristics

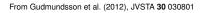
Non-reactive HiPIMS

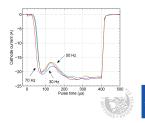


HiPIMS - Voltage - Current - time

- In non-reactive discharge the current waveform shows an initial pressure dependent peak that is followed by a second phase that is power and material dependent
- The initial phase has a contribution from the working gas ions, whereas the later phase has a strong contribution from self-sputtering at high voltage



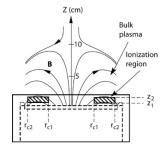




Ionization region model studies of non-reactive HiPIMS



- The ionization region model (IRM) was developed to improve the understanding of the plasma behaviour during a HiPIMS pulse and the afterglow
- The main feature of the model is that an ionization region (IR) is defined next to the race track
- The IR is defined as an annular cylinder with outer radii r_{c2} , inner radii r_{c1} and length $L = z_2 z_1$, extends from z_1 to z_2 axially away from the target



The definition of the volume covered by the IRM From Raadu et al. (2011), PSST **20** 065007



- The temporal development is defined by a set of ordinary differential equations giving the first time derivatives of
 - the electron energy
 - the particle densities for all the particles
- The species assumed in the non-reactive-IRM are
 - $\bullet\,$ cold electrons $e^{C},$ hot electrons e^{H}
 - argon atoms Ar(3s²3p⁶), warm argon atoms in the ground state Ar^W, hot argon atoms in the ground state Ar^H, Ar^m (1s₅ and 1s₃) (11.6 eV), argon ions Ar⁺ (15.76 eV)
 - titanium atoms Ti(a³F), titanium ions Ti⁺ (6.83 eV), doubly ionized titanium ions Ti²⁺ (13.58 eV)
 - aluminium atoms Al(²P_{1/2}), aluminium ions Al⁺ (5.99 eV), doubly ionized aluminium ions Al²⁺ (18.8 eV)

Detailed model description is given in Huo et al. (2017), JPD submitted 2017

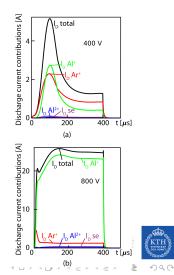
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- A non-reactive discharge with Al target
- When the discharge is operated at 400 V the contributions of Al⁺ and Ar⁺-ions to the discharge current are very similar
- At 800 V Al⁺-ions dominate the discharge current (self-sputtering) while the contribution of Ar⁺ is below 10 % except at the initiation of the pulse

From Huo et al. (2017), JPD submitted 2017

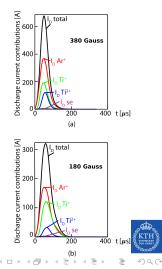
Experimental data from Anders et al. (2007) JAP 102 113303



- A non-reactive discharge with Ti target
- The contributions to the discharge current for two cases, weak (180 Gauss) and strong (380 Gauss) magnetic field, at 75 Hz pulse frequency
- Stronger magnetic field leads to a higher discharge current
- Higher magnetic field strength leads to higher relative contribution of Ti²⁺ while it lowers the relative contribution of Ti⁺

From Huo et al. (2017), JPD submitted 2017

Experimental data from Bradley et al. (2015) JPD 48 215202



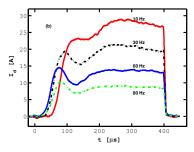
Voltage - Current - Time characteristics

Reactive HiPIMS



HiPIMS - Voltage - Current - time

- During reactive sputtering, a reactive gas is added to the inert working gas
- The current waveform in the reactive Ar/N₂ HiPIMS discharge with Ti target is highly dependent on the pulse repetition frequency
- N₂ addition changes the plasma composition and the target condition can also change due to the formation of a compound on its surface

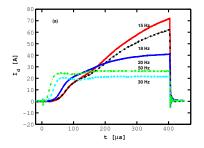


After Magnus et al. (2011) JAP 110 083306



HiPIMS - Voltage - Current - time

- Similarly for the Ar/O₂ discharge, the current waveform is highly dependent on the repetition frequency and applied voltage which is linked to oxide formation on the target
- The current is found to increase significantly as the frequency is lowered



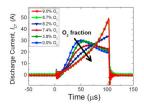
After Magnus et al. (2012), JVSTA 30 050601

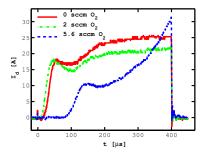
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HiPIMS - Voltage - Current - time

 As the oxygen flow is increased a transition to oxide mode is observed



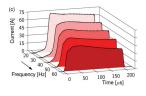


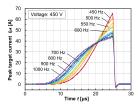
The current waveforms for an $\mbox{Ar/O}_2$ discharge with a V target where the oxygen flow rate is varied

From Aijaz et al. (2016) Solar Energy Materials and Solar Cells **149** 137

The current waveforms for an Ar/O₂ discharge with a Ti target where the oxygen flow rate is varied – 600 V, 50 Hz and 0.6 Pa From Gudmundsson et al. (2013), ISSP 2013, p. 192 Gudmundsson (2016) Plasma Phys. Contr. Fus. **58** 014002

HiPIMS - Voltage - Current - time





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- Similar behaviour has been reported for various target and reactive gas combinations
 - The current increases with decreased repetition frequency
 - The current waveform maintains its shape for Ar/O₂ discharge with Nb target

From Hála et al. (2012), JPD 45 055204

 The current waveform becomes distinctly triangular for Ar/N₂ discharge with Hf target

From Shimizu et al. (2016), JPD 49 065202



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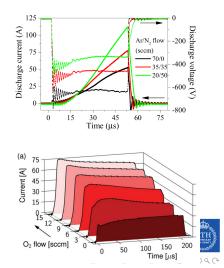
HiPIMS - Voltage - Current - time

- The current increases with increased partial pressure of the reactive gas
 - The current waveform becomes distinctly triangular for Ar/N₂ discharge with Al target

From Moreira et al. (2015), JVSTA 33 021518

 The current waveform maintains its shape for Ar/O₂ discharge with Nb target

From Hála et al. (2012), JPD 45 055204



Ionization region model studies of reactive HiPIMS



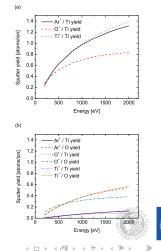
• The species assumed in the reactive-IRM are

- $\bullet\,$ cold electrons $e^{\rm C},$ hot electrons $e^{\rm H}$
- argon atoms Ar(3s²3p⁶), warm argon atoms in the ground state Ar^W, hot argon atoms in the ground state Ar^H, Ar^m (1s₅ and 1s₃) (11.6 eV), argon ions Ar⁺ (15.76 eV)
- titanium atoms Ti(a³F), titanium ions Ti⁺ (6.83 eV), doubly ionized titanium ions Ti²⁺ (13.58 eV)
- oxygen molecule in the ground state $O_2(X^3\Sigma_g^-)$, the metastable oxygen molecules $O_2(a^1\Delta_g)$ (0.98 eV) and $O_2(b^1\Sigma_g)$ (1.627 eV), the oxygen atom in the ground state $O(^3P)$, the metastable oxygen atom $O(^1D)$ (1.96 eV), the positive ions O_2^+ (12.61 eV) and O^+ (13.62 eV), and the negative ion O^-



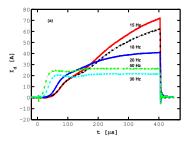
- The sputter yield for the various bombarding ions was calculated by TRIDYN for
 - Metal mode Ti target
 - Poisoned mode TiO₂ target
- The yields correspond to the extreme cases of either clean Ti surface and a surface completely oxidized (TiO₂ surface)
- The sputter yield is much lower for poisoned target

The sputter yield data is from Tomas Kubart, Uppsala University





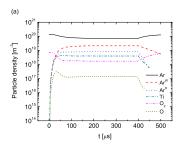
- The model is applied to explore Ar/O₂ discharge with Ti target in both metal mode and oxide (poisoned) mode
- The IRM is a semi-empirical model in the sense that it uses a measured discharge voltage and current waveforms as a main input parameter
- For this study we use the measured curve for Ar/O₂ with Ti target at 50 Hz for metal mode and at 15 Hz for poisoned mode



After Magnus et al. (2012), JVSTA 30 050601



- The gas rarefaction is observed for the argon atoms but is more significant for the O₂ molecule
- The density of Ti atoms is higher than the O₂ density
- The atomic oxygen density of is over one order of magnitude lower than the molecular oxygen density – the dissociation fraction is low



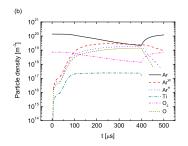
The temporal evolution of the neutral species with 5 % oxygen partial flow rate for Ar/O₂ discharge with Ti target in metal mode.

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Gudmundsson et al. (2016), PSST, 25(6) 065004 KTH

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- Gas rarefaction is observed for both argon atoms and O₂ molecules
- The density of Ti atoms is lower than both the O₂ density and atomic oxygen density
- The atomic oxygen density is higher than the O₂ density towards the end of the pulse



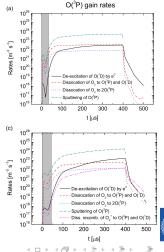
The temporal evolution of the neutral species with **5** % oxygen partial flow rate for Ar/O₂ discharge with Ti target in **poisoned mode**. Gudmundsson et al. (2016), PSST, 25(6) 065002

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- The increase in the atomic oxygen in the ground state is due to:
 - sputtering of O(³P) from the partially to fully oxidized target (dominates)
 - electron impact de-excitation of O(¹D)
 - electron impact dissociation of the O₂ ground state molecule

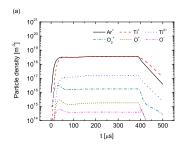
The temporal evolution of the neutral species with **5** % oxygen partial flow rate for Ar/O₂ discharge with Ti target in transition mode and poisoned mode.

Lundin et al. (2017), JAP, 121(17) 171917



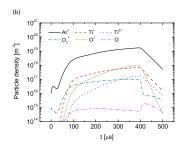
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- Ar⁺ and Ti⁺-ions dominate the discharge
- Ti²⁺-ions follow by roughly an order of magnitude lower density
- The O₂⁺ and O⁺-ion density is much lower



The temporal evolution of the neutral species with 5 % oxygen partial flow rate for Ar/O₂ discharge with Ti target in metal mode. Gudmundsson et al. (2016), PSST, 25(6) 06500

- Ar⁺-ions dominate the discharge
- Ti⁺, O⁺, have very similar density, but the temporal variation is different, and the O₂⁺ density is slightly lower
- The Ti²⁺-ion density increases fast with time and overcomes the O₂⁺ density towards the end of the pulse



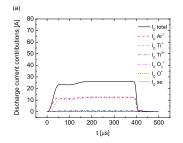
The temporal evolution of the neutral species with **5** % oxygen partial flow rate for Ar/O₂ discharge with Ti target in poisoned mode. Gudmundsson et al. (2016), PSST, 25(6) 06500

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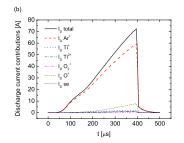
Ionization region model studies of reactive HiPIMS

 Ar⁺ and Ti⁺-ions contribute most significantly to the discharge current at the cathode target surface – almost equal contribution



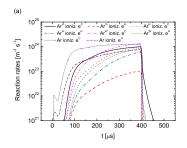
The temporal evolution of the neutral species with 5 % oxygen partial flow rate for Ar/O₂ discharge with Ti target in metal mode. Gudmundsson et al. (2016), PSST, 25(6) 06500

- Ar⁺ contribute most significantly to the discharge current – almost solely – at the cathode target surface
- The contribution of secondary electron emission is very small



The temporal evolution of the neutral species with 5 % oxygen partial flow rate for Ar/O₂ discharge with Ti target in poisoned mode. Gudmundsson et al. (2016), PSST, 25(6) 06500

- Recycling of atoms coming from the target and then ionized are required for the current generation in both modes of operation
- In the metal mode self-sputter recycling dominates and in the poisoned mode working gas recycling dominates
- The dominating type of recycling determines the discharge current waveform



The temporal variations of the reaction rates for electron impact ionization of the argon atoms (ground state plus metastable) in poisoned mode.

Gudmundsson et al. (2016), PSST, 25(6) 06

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- In the metal mode sheath energization was found to be only 10 %
 - same range as the results reported earlier for an AI target Huo et al. (2013), PSST 22(4) (2013) 045005

the dominating electron heating mechanism is Ohmic heating

- For the poisoned mode the sheath energization was 30 %, with a rising trend, at the end of the pulse
- This is due to the secondary electron emission
 - In the poisoned mode essentially all the ions (mainly Ar⁺, but also O⁺ and Ti²⁺ towards the end of the pulse) contribute to the secondary electron emission
 - In the metal mode only half of the ions contribute to the secondary electron emission (Ar⁺) while the other half does not contribute at all ($\gamma_{Ti^+} = 0.0$)



Summary



Summary

- An ionization region model was used to explore the plasma composition during the high power pulse
- Comparison was made between the metal mode and the poisoned mode
 - In metal mode Ar⁺ and Ti⁺-ions dominate the discharge and are of the same order of magnitude
 - In poisoned mode Ar⁺-ions dominate the discharge and two orders of magnitude lower, Ti⁺, O⁺, have very similar density, with the O₂⁺ density slightly lower
 - In the metal mode Ar⁺ and Ti⁺-ions contribute most significantly to the discharge current while in poisoned mode Ar⁺ dominate
- In the metal mode self-sputter recycling dominates and in the poisoned mode working gas recycling dominates – the dominating type of recycling determines the discharge current waveform



The slides can be downloaded at

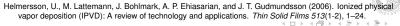
http://langmuir.raunvis.hi.is/~tumi/ranns.html

- We got help with the sputtering yields from
 - Dr. Tomas Kubart, Uppsala University, Uppsala, Sweden
- and the project is funded by
 - Icelandic Research Fund Grant No. 130029
 - Swedish Government Agency for Innovation Systems (VINNOVA) contract no. 2014-04876,



References

- Aijaz, A., Y.-X. Ji, J. Montero, G. A. Niklasson, C. G. Granqvist, and T. Kubart (2016). Low-temperature synthesis of thermochronic vanadium dioxide thin films by reactive high power impulse magnetron sputtering. *Solar Energy Materials and Solar Cells* 149, 137–144.
- Aiempanakit, M., A. Aijaz, D. Lundin, U. Helmersson, and T. Kubart (2013). Understanding the discharge current behavior in reactive high power impulse magnetron sputtering of oxides. *Journal of Applied Physics* 113(13), 133302.
- Anders, A., J. Andersson, and A. Ehiasarian (2007). High power impulse magnetron sputtering: Current-voltage-time characteristics indicate the onset of sustained self-sputtering. *Journal of Applied Physics 102*(11), 113303.
- Bradley, J. W., A. Mishra, and P. J. Kelly (2015). The effect of changing the magnetic field strength on HiPIMS deposition rates. *Journal of Physics D: Applied Physics* 48(21), 215202.
- Gudmundsson, J. T. (2016). On reactive high power impulse magnetron sputtering. Plasma Physics and Controlled Fusion 58(1), 014002.
- Gudmundsson, J. T., N. Brenning, D. Lundin, and U. Helmersson (2012). The high power impulse magnetron sputtering discharge. *Journal of Vacuum Science and Technology A* 30(3), 030801.
- Gudmundsson, J. T., F. Magnus, T. K. Tryggvason, S. Shayestehaminzadeh, O. B. Sveinsson, and S. Olafsson (2013). Reactive high power impulse magnetron sputtering. In *Proceedings of the XII International Symposium* on Sputtering and Plasma Processes (ISSP 2013), pp. 192–194.
- Gudmundsson, J. T., D. Lundin, N. Brenning, M. A. Raadu, C. Huo, and T. M. Minea (2016). An ionization region model of the reactive Ar/O₂ high power impulse magnetron sputtering discharge. *Plasma Sources Science and Technology 25*(6), 065004.
- Hála, M., J. Čapek, O. Zabeida, J. E. Klemberg-Sapieha, and L. Martinu (2012). Hysteresis free deposition of niobium oxide films by HiPIMS using different pulse management strategies. *Journal of Physics D: Applied Physics* 45(5), 055204.





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References

- Huo, C., D. Lundin, M. A. Raadu, A. Anders, J. T. Gudmundsson, and N. Brenning (2013). On sheath energization and ohmic heating in sputtering magnetrons. *Plasma Sources Science and Technology 22*(4), 045005.
- Huo, C., D. Lundin, J. T. Gudmundsson, M. A. Raadu, J. W. Bradley, and N. Brenning (submitted 2017). Particle-balance models for pulsed sputtering magnetrons. *Journal of Physics D: Applied Physics*.
- Lundin, D., J. T. Gudmundsson, N. Brenning, M. A. Raadu, and T. M. Minea (2017). A study of the oxygen dynamics in a reactive Ar/O₂ high power impulse magnetron sputtering discharge using an ionization region model. *Journal of Applied Physics* 121(17), 171917.
- Magnus, F., A. S. Ingason, S. Olafsson, and J. T. Gudmundsson (2012). Nucleation and resistivity of ultrathin TiN films grown by high power impulse magnetron sputtering. *IEEE Electron Device Letters* 33(7), 1045 – 1047.
- Magnus, F., O. B. Sveinsson, S. Olafsson, and J. T. Gudmundsson (2011). Current-voltage-time characteristics of the reactive Ar/N₂ high power impulse magnetron sputtering discharge. *Journal of Applied Physics* 110(8), 083306.
- Magnus, F., T. K. Tryggvason, S. Olafsson, and J. T. Gudmundsson (2012). Current-voltage-time characteristics of the reactive Ar/O₂ high power impulse magnetron sputtering discharge. *Journal of Vacuum Science and Technology A* 30(5), 050601.
- Moreira, M. A., T. Törndahl, I. Katardjiev, and T. Kubart (2015). Deposition of highly textured AIN thin films by reactive high power impulse magnetron sputtering. *Journal of Vacuum Science and Technology A* 33(2), 021518.
- Raadu, M. A., I. Axnäs, J. T. Gudmundsson, C. Huo, and N. Brenning (2011). An ionization region model for high power impulse magnetron sputtering discharges. *Plasma Sources Science and Technology* 20(6), 065007.
- Shimizu, T., M. Villamayor, D. Lundin, and U. Helmersson (2015). Process stabilization by peak current regulation in reactive high-power impulse magnetron sputtering of hafnium nitride. *Journal of Physics D: Applied Physics* 49(6), 065202.



Toneli, D. A., R. S. Pessoa, M. Roberto, and J. T. Gudmundsson (2015). On the formation and annihilation of the singlet molecular metastables in an oxygen discharge. *Journal of Physics D: Applied Physics 48*(32), 325202.

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