

# The current waveform in reactive high power impulse magnetron sputtering

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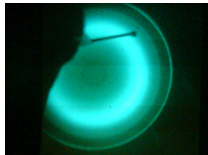
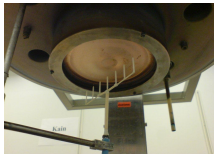
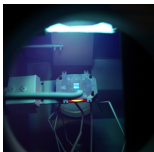
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4<sup>th</sup> Magnetron, Ion processing & Arc Technologies European Conference,  
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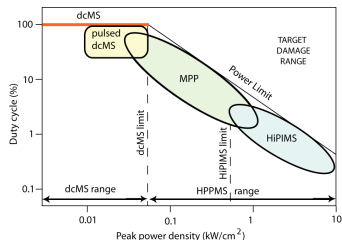
# Introduction



- Magnetron sputtering has been the workhorse of plasma based sputtering methods for over four decades
- For many applications a high degree of ionization of the sputtered vapor is desired
  - controlled ion bombardment of the growing film – controlled by a negative bias applied to the substrate
  - collimation – enhanced step coverage
- Common to all highly ionized magnetron sputtering techniques is a very high density plasma

# Introduction

- In a conventional dc magnetron discharge the power density is limited by the thermal load on the target
- High power pulsed magnetron sputtering (HPPMS)
- In a HiPIMS discharge a high power pulse is supplied for a short period
  - low frequency
  - low duty cycle
  - low average power

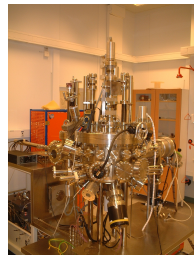


From Gudmundsson et al. (2012), JVSTA **30** 030801

- Power density limits  
 $p_t = 0.05 \text{ kW/cm}^2$  dcMS limit  
 $p_t = 0.5 \text{ kW/cm}^2$  HiPIMS limit

# Introduction

- Reactive sputtering, where metal targets are sputtered in a reactive gas atmosphere to deposit compound materials is of utmost importance in various technologies
- In reactive sputtering processes a reactive gas  $O_2$ ,  $N_2$ , or  $CH_4$  etc. is mixed to the noble working gas for oxide, nitride, or carbide deposition
- The high electron density in the HiPIMS discharge is expected to enhance the dissociation of the molecular gas

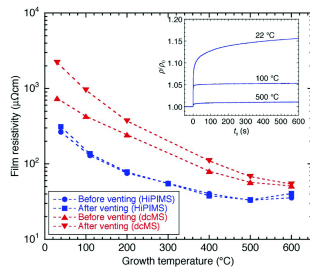


# Reactive HiPIMS - Applications



## Application – Film Resistivity

- TiN as diffusion barriers for interconnects
- HiPIMS deposited films have significantly lower resistivity than dcMS deposited films on  $\text{SiO}_2$  at all growth temperatures due to reduced grain boundary scattering
- Thus, ultrathin continuous TiN films with superior electrical characteristics and high resistance towards oxidation can be obtained with HiPIMS at reduced temperatures

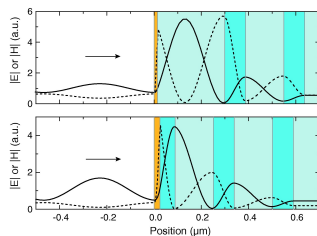


From Magnus et al. (2012) IEEE EDL **33** 1045

## Application – Bragg mirror

- Multilayer structures containing a high-contrast ( $\text{TiO}_2/\text{SiO}_2$ ) Bragg mirror fabricated on fused-silica substrates
  - reactive HiPIMS  $\text{TiO}_2$  (88 nm)
  - reactive dcMS  $\text{SiO}_2$  (163 nm)
  - capped with semitransparent gold
- Rutile  $\text{TiO}_2$  ( $n = 2.59$ ) and  $\text{SiO}_2$  ( $n = 1.45$ ) provide a large index contrast
- Smooth rutile  $\text{TiO}_2$  films can be obtained by HiPIMS at relatively low growth temperatures, without post-annealing

Agnarsson et al. (2013) TSF 545 445



From Leosson et al. (2012) Opt. Lett. 37 4026

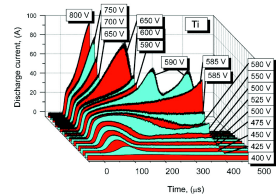
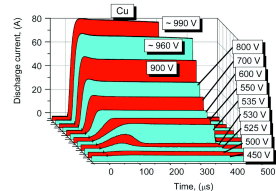
# Reactive HiPIMS - Voltage - Current - Time characteristics





## HiPIMS - Voltage - Current - time

- In **non-reactive** discharge the current waveform shows an initial pressure dependent peak that is followed by a second phase that is power and material dependent
- The initial phase is dominated by gas ions, whereas the later phase has a strong contribution from self-sputtering
- For some materials, the discharge switches into a mode of **sustained self-sputtering**



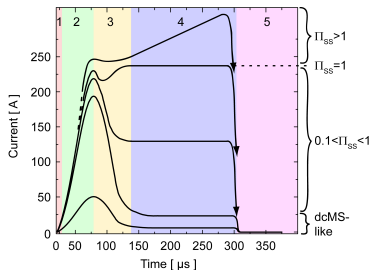
From Anders et al. (2007),

JAP 102 113303 and JAP 103 039901



## HiPIMS - Voltage - Current - time

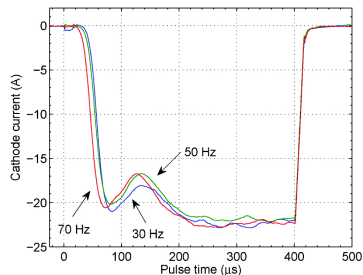
- A schematic illustration of the discharge current assuming square shaped voltage pulses
- The current is generally characterized by an initial peak followed by a more or less stable current plateau (bottom current curves)
- In other cases it shows an initial peak followed by a second increase of the discharge current (top current curves)
- The non-reactive case is well understood



From Gudmundsson et al. (2012), JVSTA **30** 030801

## HiPIMS - Voltage - Current - time

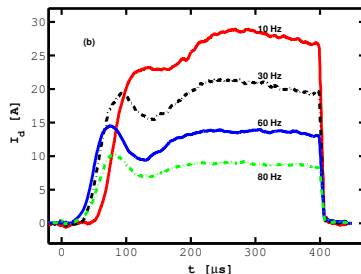
- Ar discharge with Ti target
- The initial peak in current results large flux of atoms from the target
- Collisions of the sputtered atoms with the working gas result in heating and expansion of the working gas – **rarefaction**
- A significant fraction of the sputtered atoms experience electron impact ionization (the ionization mean free path  $\sim 1$  cm) and are attracted back to the target to participate in the sputtering process – **self-sputtering**



From Magnus et al. (2011) JAP **110** 083306

## HiPIMS - Voltage - Current - time

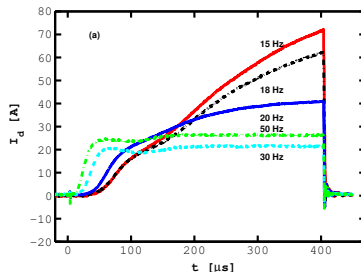
- During reactive sputtering, a reactive gas is added to the inert working gas
- The current waveform in the reactive Ar/N<sub>2</sub> HiPIMS discharge is highly dependent on the pulse repetition frequency, unlike for pure Ar
- N<sub>2</sub> addition changes the plasma composition and the target condition can also change due to the formation of a compound on its surface



After Magnus et al. (2011) JAP **110** 083306

## HiPIMS - Voltage - Current - time

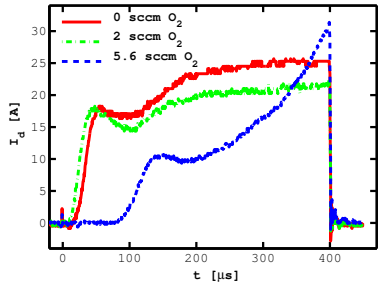
- Similarly for the Ar/O<sub>2</sub> discharge, the current waveform is highly dependent on the repetition frequency and applied voltage which is linked to oxide formation on the target
- The current is found to increase significantly as the frequency is lowered



After Magnus et al. (2012), JVSTA **30** 050601

## HiPIMS - Voltage - Current - time

- As the oxygen flow is increased a transition to oxide mode is observed – The delay in the onset of the current increases, the initial current peak is lowered and a transition to a self-sputtering runaway occurs
- It has been confirmed that in the oxide mode, the discharge is dominated by  $O^+$ -ions, due to oxygen atoms sputtered off the target surface



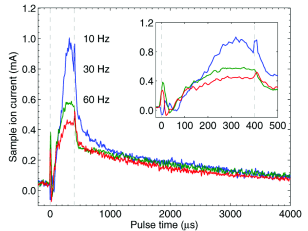
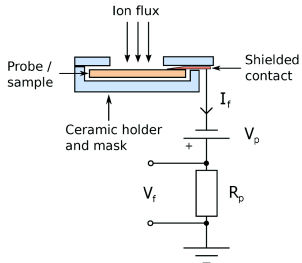
The current waveforms for an Ar/O<sub>2</sub> discharge with a Ti target where the oxygen flow rate is varied – 600 V, 50 Hz and 0.6 Pa

From Gudmundsson et al. (2013), ISSP 2013, p. 192

Gudmundsson (2016) Plasma Phys. Contr. Fus. **58** 014002

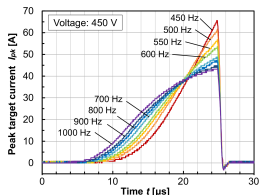
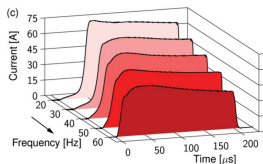
# HiPIMS - Voltage - Current - time

- The observed changes in the discharge current are reflected in the flux of ions impinging on the substrate



From Magnus et al. (2011), JAP **110** 083306

# HiPIMS - Voltage - Current - time



- Similar behaviour has been reported for various target and reactive gas combinations
  - The current increases with decreased repetition frequency
  - The current waveform maintains its shape for Ar/O<sub>2</sub> discharge with Nb target
- The current waveform becomes distinctly triangular for Ar/N<sub>2</sub> discharge with Hf target

From Hála et al. (2012), JPD **45** 055204

From Shimizu et al. (2015), arXiv:1509.07002

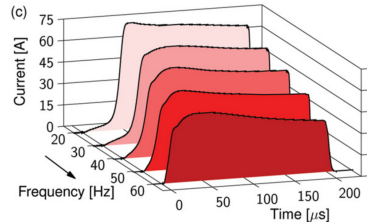
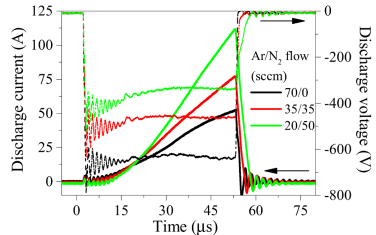


## HiPIMS - Voltage - Current - time

- The current increases with increased partial pressure of the reactive gas
  - The current waveform becomes distinctly triangular for Ar/N<sub>2</sub> discharge with Al target
- The current waveform maintains its shape for Ar/O<sub>2</sub> discharge with Nb target

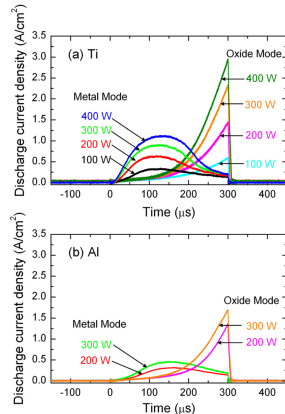
From Moreira et al. (2015), JVSTA **33** 021518

From Hála et al. (2012), JPD **45** 055204



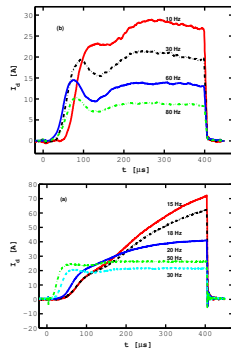
# HiPIMS - Voltage - Current - time

- The shape of the current waveform depends on the mode of operation
  - The current waveform becomes distinctly triangular for Ar/O<sub>2</sub> discharge with both Al and Ti target



### *HiPIMS - Voltage - Current - time*

- HiPIMS differs significantly from dcMS, due to the fact that self-sputtering quickly becomes dominant and the working gas ions (mostly  $\text{Ar}^+$  and  $\text{N}_2^+$  or  $\text{O}_2^+$ ) are depleted from the area in front of the target, due to rarefaction
- The secondary electron emission yield is governed by the composition of the target (Ti or TiN or  $\text{TiO}_2$ ) and the type of ions that are bombarding it

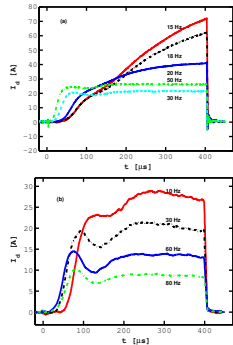


After Magnus et al. (2011), JAP **110** 083306

and Magnus et al. (2012), JVSTA 30 05060

## HiPIMS - Voltage - Current - time

- $\gamma_{\text{see}}$  is practically zero for singly charged metal ions impacting a target of the same metal
- $\gamma_{\text{see}}$  will be higher for self sputtering from a TiN or TiO<sub>2</sub> target, where N<sup>+</sup>-ions or O<sup>+</sup>-ions are also present, than for self-sputtering from a Ti target, where multiply charged Ti ions are needed to create secondary electrons

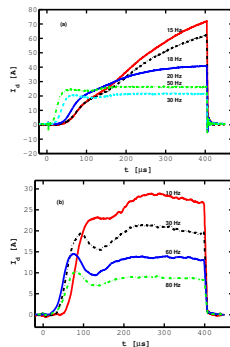


After Magnus et al. (2011), JAP **110** 083306

and Magnus et al. (2012), JVSTA **30** 050601

## HiPIMS - Voltage - Current - time

- At high frequencies, nitride or oxide is not able to form between pulses, and self-sputtering by  $\text{Ti}^+$ -ions (singly and multiply charged) from a Ti target is the dominant process
- At low frequency, the long off-time results in a nitride or oxide layer being formed on the target surface and self-sputtering by  $\text{Ti}^+$ - and  $\text{N}^+$ -ions or  $\text{O}^+$ -ions from  $\text{TiN}$  or  $\text{TiO}_2$  takes place



From Magnus et al. (2011), JAP **110** 083306

and Magnus et al. (2012), JVSTA **30** 050601

Gudmundsson (2016) PPCF **58** 014002

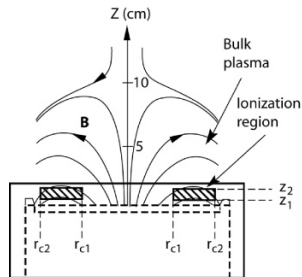


# Ionization region model studies of reactive HiPIMS



# *Ionization region model studies of reactive HiPIMS*

- The ionization region model (IRM) was developed to improve the understanding of the plasma behaviour during a HiPIMS pulse and the afterglow
- The main feature of the model is that an ionization region (IR) is defined next to the race track
- The IR is defined as an annular cylinder with outer radii  $r_{c2}$ , inner radii  $r_{c1}$  and length  $L = z_2 - z_1$ , extends from  $z_1$  to  $z_2$  axially away from the target

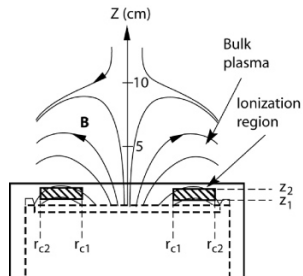


The definition of the volume covered by the IRM

From Raadu et al. (2011), PSST **20** 065007

## ***Ionization region model studies of reactive HiPIMS***

- Geometrical effects are included indirectly as loss and gain rates across the boundaries of this annular cylinder to the target and the bulk plasma
- The temporal development is defined by a set of ordinary differential equations giving the first time derivatives of the electron energy and the particle densities for all the particles
- The electron density is found assuming quasi-neutrality of the plasma



The definition of the volume covered by the IRM

From Raadu et al. (2011), PSST **20** 065007

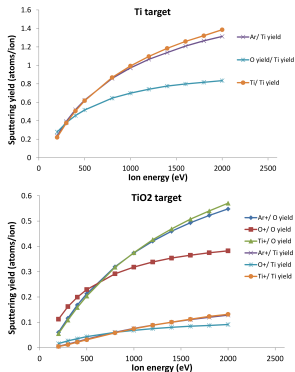


# ***Ionization region model studies of reactive HiPIMS***

- The species assumed in the IRM are
  - electrons
  - argon atoms  $\text{Ar}(3s^23p^6)$ , warm argon atoms in the ground state  $\text{Ar}^{\text{w}}$ , hot argon atoms in the ground state  $\text{Ar}^{\text{H}}$ ,  $\text{Ar}^{\text{m}}$  ( $1s_5$  and  $1s_3$ ) (11.6 eV), argon ions  $\text{Ar}^+$  (15.76 eV), doubly ionized argon ions  $\text{Ar}^{2+}$  (27.63 eV)
  - titanium atoms  $\text{Ti}(a^3F)$ , titanium ions  $\text{Ti}^+$  (6.83 eV), doubly ionized titanium ions  $\text{Ti}^{2+}$  (13.58 eV)
  - oxygen molecule in the ground state  $\text{O}_2(X^3\Sigma_g^-)$ , the metastable oxygen molecules  $\text{O}_2(a^1\Delta_g)$  (0.98 eV) and  $\text{O}_2(b^1\Sigma_g)$  (1.627 eV), the oxygen atom in the ground state  $\text{O}(^3P)$ , the metastable oxygen atom  $\text{O}(^1D)$  (1.96 eV), the positive ions  $\text{O}_2^+$  (12.61 eV) and  $\text{O}^+$  (13.62 eV), and the negative ion  $\text{O}^-$

# *Ionization region model studies of reactive HiPIMS*

- The sputter yield for the various bombarding ions was calculated by TRIDYN for
  - Metal mode – Ti target
  - Poisoned mode – TiO<sub>2</sub> target
- The yields correspond to the extreme cases of either clean Ti surface and a surface completely oxidized (TiO<sub>2</sub> surface)

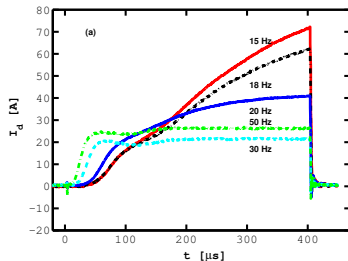


The sputter yield data is from Tomas Kubart,  
Uppsala University



## *Ionization region model studies of reactive HiPIMS*

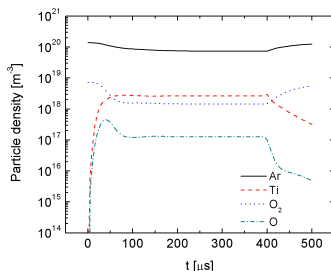
- The model is applied to explore Ar/O<sub>2</sub> discharge with Ti target in both metal mode and oxide (poisoned) mode
- The IRM is a semi-empirical model in the sense that it uses a measured discharge current waveform as a main input parameter
- For this study we use the measured curve for Ar/O<sub>2</sub> with Ti target at 50 Hz for metal mode and at 15 Hz for poisoned mode



After Magnus et al. (2012), JVSTA **30** 050601

# ***Ionization region model studies of reactive HiPIMS***

- The gas rarefaction is observed for the argon atoms but is more significant for the  $O_2$  molecule
- The density of Ti atoms is higher than the  $O_2$  density
- The atomic oxygen density of is over one order of magnitude lower than the molecular oxygen density – the dissociation fraction is low

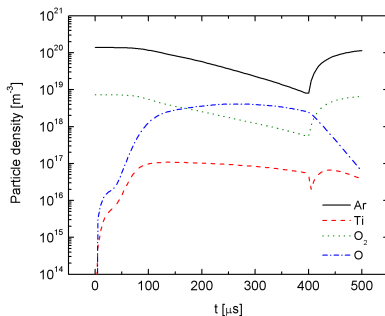


The temporal evolution of the neutral species with 5 % oxygen partial flow rate for Ar/ $O_2$  discharge with Ti target in **metal mode**.



# ***Ionization region model studies of reactive HiPIMS***

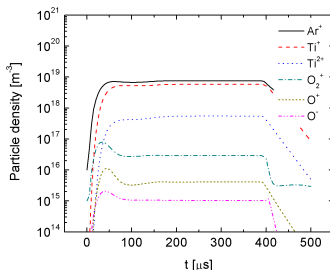
- Gas rarefaction is observed for both argon atoms and O<sub>2</sub> molecules
- The density of Ti atoms is lower than both the O<sub>2</sub> density and atomic oxygen density
- The atomic oxygen density is higher than the O<sub>2</sub> density in towards the end of the pulse



The temporal evolution of the neutral species with 5 % oxygen partial flow rate for Ar/O<sub>2</sub> discharge with Ti target in **poisoned mode**.

# ***Ionization region model studies of reactive HiPIMS***

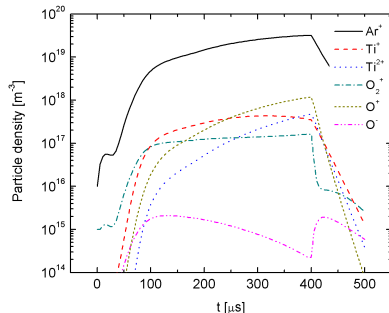
- $\text{Ar}^+$  and  $\text{Ti}^+$ -ions dominate the discharge
- $\text{Ti}^{2+}$ -ions follow by roughly an order of magnitude lower density
- The  $\text{O}_2^+$  and  $\text{O}^+$ -ion density is much lower



The temporal evolution of the neutral species with **5 % oxygen partial flow rate** for  $\text{Ar}/\text{O}_2$  discharge with Ti target in **metal mode**.

# ***Ionization region model studies of reactive HiPIMS***

- $\text{Ar}^+$ -ions dominate the discharge
- $\text{Ti}^+$ ,  $\text{O}^+$ , have very similar density, but the temporal variation is different, and the  $\text{O}_2^+$  density is slightly lower
- The  $\text{Ti}^{2+}$ -ion density increases fast with time and overcomes the  $\text{Ti}^+$  density towards the end of the pulse

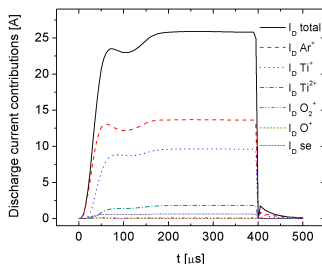


The temporal evolution of the neutral species with 5 % oxygen partial flow rate for Ar/ $\text{O}_2$  discharge with Ti target in **poisoned mode**.



# ***Ionization region model studies of reactive HiPIMS***

- $\text{Ar}^+$  and  $\text{Ti}^+$ -ions contribute most significantly to the discharge current

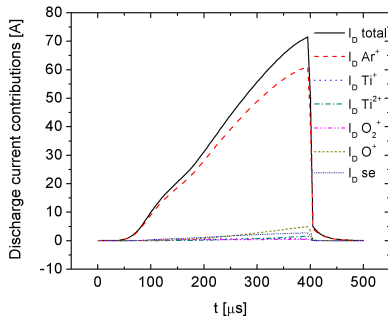


The temporal evolution of the neutral species with 5 % oxygen partial flow rate for Ar/O<sub>2</sub> discharge with Ti target in **metal mode**.



# ***Ionization region model studies of reactive HiPIMS***

- $\text{Ar}^+$  contribute most significantly to the discharge current – almost solely
- The contribution of secondary electron emission is very small



The temporal evolution of the neutral species with 5 % oxygen partial flow rate for  $\text{Ar}/\text{O}_2$  discharge with Ti target in **poisoned mode**.



# Summary



## Summary

- The current-voltage-time waveforms in a reactive discharge exhibit similar general characteristics as the non-reactive case in some cases
  - the current rises to a peak, then decays because of rarefaction before rising to a self-sputtering dominated phase
  - in other cases the current develops a triangular shape as repetition frequency is lowered or the partial pressure of the reactive gas is increased
- At low repetition frequency, the long off-time results in a nitride or oxide layer being formed on the target surface and an increase in the discharge current



## Summary

- An ionization region model was used to explore the plasma composition during the high power pulse
- Comparison was made between the metal mode and the poisoned mode
  - In metal mode  $\text{Ar}^+$  and  $\text{Ti}^+$ -ions dominate the discharge and are of the same order of magnitude
  - In poisoned mode  $\text{Ar}^+$ -ions dominate the discharge and two orders of magnitude lower,  $\text{Ti}^+$ ,  $\text{O}^+$ , have very similar density, with the  $\text{O}_2^+$  density slightly lower
  - In the metal mode  $\text{Ar}^+$  and  $\text{Ti}^+$ -ions contribute most significantly to the discharge current while in poisoned mode  $\text{Ar}^+$  dominate
- $\text{Ar}^+$ -ions are responsible for the increase in the discharge current in poisoned mode



The slides can be downloaded at

<http://langmuir.raunvis.hi.is/~tumi/hipims.html>

- The experimental work was made in collaboration with
  - Dr. Fridrik Magnus, Uppsala University, Uppsala, Sweden
  - Tryggvi K. Tryggvason, University of Iceland
- The modeling work is in collaboration with
  - Dr. Daniel Lundin, Université Paris-Sud, Orsay, France
  - Prof. Nils Brenning, KTH Royal Institute of Technology, Stockholm, Sweden
  - Prof. Tiberu Minea, Université Paris-Sud, Orsay, France
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